



From linear potential theory to the inverse mean curvature flow

Monotonicity formulas

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Geometric Measure Theory on Metric Spaces with Applications to Physics and Geometry

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Result

Regularity/geometric interpretation of the level sets of p-harmonic functions.

Aims

- link monotonicity formulas along the inverse mean curvature flow to monotonicity formulas in (non-)linear potential theory,
- o justify formal computations,
- o geometric inequalities.

Setting

 (M^n,g) complete noncompact Riemannian n-manifold, $n \ge 3$.

- \circ **Setting 1**: nonnegative Ricci curvature $\mathrm{Ric} \geq 0$;
- \circ **Setting 2**: Riemannian 3-manifolds, nonnegative scalar curvature $R \ge 0$.

IMCF and harmonic functions	1
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Limitations of the smooth setting	3
GMT meets PDE	4

1 IMCF and harmonic functions

2

Definition - Inverse Mean Curvature Flow

Let N be an hypersurface in (M^n,g) a n-Riemannian smooth manifold, and $F_0:N\to M$ a smooth immersion. A classical solution to the Inverse Mean Curvature flow is a family of smooth immersions $F:[0,T)\times N\to M$ satisfying

$$\begin{cases} \partial_t F(p,t) = \frac{\nu(F(t,p))}{\mathrm{H}(F(t,p))} \\ F(0,p) = F_0(p) \end{cases}$$

(IMCF)

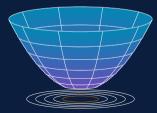
where H is the mean curvature and ν the outward normal vector.

Expanding spheres

Let $N=\mathbb{S}^{n-1}$, $M=\mathbb{R}^n$, $F_0(N)=\mathbb{S}^{n-1}_{R_0}.$ Then, IMCF reduces to the following ODE

$$\begin{cases} R'(t) = \frac{1}{n-1}R(t) \\ R(0) = R_0 \end{cases}$$

thus,
$$R(t) = R_0 e^{\frac{1}{(n-1)t}}$$
.



MONOTONE QUANTITIES

Setting I: (M,g) complete, noncompact Riemannian 3-manifold with nonnegative Ricci curvature $\operatorname{Ric} > 0$, Σ hypersurface in M.

Let Σ_{\star} be the IMCF from Σ_{\star} . We compute the evolution of the area, the volume and the mean curvature:

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}t}|\Sigma_t| &= |\Sigma_t|, \qquad \frac{\mathrm{d}}{\mathrm{d}t}\mathrm{Vol}(\Sigma_t) = \int_{\Sigma_t} \frac{1}{\mathrm{H}}\,\mathrm{d}\sigma, \\ \frac{\mathrm{d}}{\mathrm{d}t}\mathrm{H} &= -\Delta \left(\frac{1}{\mathrm{H}}\right) - \frac{|\mathrm{h}|^2}{\mathrm{H}} - \frac{\mathrm{Ric}(\nu,\nu)}{\mathrm{H}}. \end{split}$$

Then, the evolution of the Willmore functional reads

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}t} \int_{\Sigma_t} \mathrm{H}^2 \, \mathrm{d}\sigma &= \int_{\Sigma_t} -2 \mathrm{H}\Delta(1/\mathrm{H}) -2 \mathrm{H}\Delta(1/\mathrm{H}) -2 |\mathbf{h}|^2 -2 \mathrm{Ric}(\nu,\nu) + \mathrm{H}^2 \mathrm{d}\sigma \\ &= 2 |\mathbf{h}|^2 -\mathrm{H}^2 = 2 |\mathring{\mathbf{h}}|^2 \int_{\Sigma_t} -2 \frac{|\nabla^\top \mathrm{H}|^2}{\mathrm{H}^2} -2 |\mathring{\mathbf{h}}|^2 -2 |\mathring{\mathbf{h}}|$$

where we used $0 = \int \operatorname{div}(H\nabla(1/H)) d\sigma = \int H\Delta(1/H) - |\nabla^{\top}H|^2 d\sigma$.

 $\int_{\Sigma_{-}} \mathrm{H}^2 \,\mathrm{d}\sigma$ is monotone <u>non-increasing</u>.

Setting 2: (M,g) complete, noncompact Riemannian 3-manifold with nonnegative scalar curvature $R \ge 0$, Σ hypersurface in M.

Consider the Hawking mass

$$\mathfrak{m}_H(\Sigma) \coloneqq \sqrt{\frac{|\Sigma|}{16\pi}} \left(1 - \frac{1}{16\pi} \int_{\Sigma} \mathbf{H}^2 \, \mathrm{d}\sigma \right).$$

Let Σ_t be the IMCF from Σ . We compute

$$\begin{split} \frac{\mathrm{d}}{\mathrm{d}t} & \mathfrak{m}_{H}(\Sigma_{t}) = \frac{1}{(16\pi)^{3/2}} \left[\frac{\mathrm{d}}{\mathrm{d}t} \sqrt{|\Sigma_{t}|} \left(16\pi - \int_{\Sigma_{t}} \mathrm{H}^{2} \, \mathrm{d}\sigma \right) - \sqrt{|\Sigma_{t}|} \left(\frac{\mathrm{d}}{\mathrm{d}t} \int_{\Sigma_{t}} \mathrm{H}^{2} \, \mathrm{d}\sigma \right) \right] \\ & = \frac{1}{16\pi} \sqrt{\frac{|\Sigma_{t}|}{16\pi}} \left(8\pi - \int_{\Sigma_{t}} \frac{\mathrm{H}^{2}}{2} \, \mathrm{d}\sigma + \int_{\Sigma_{t}} 2 \frac{|\nabla^{\top} \mathrm{H}|^{2}}{\mathrm{H}^{2}} + 2 |\mathring{\mathrm{h}}|^{2} + 2 \mathrm{Ric}(\nu, \nu) \, \mathrm{d}\sigma \right) \\ & = \frac{2 \mathrm{Ric}(\nu, \nu) - \mathrm{H}^{2}/2 = \mathrm{R} - \mathrm{R}^{\top} - |\mathring{\mathrm{h}}|^{2}}{2} \frac{1}{16\pi} \sqrt{\frac{|\Sigma_{t}|}{16\pi}} \left(8\pi - \int_{\Sigma_{t}} \frac{\mathrm{R}^{\top} \, \mathrm{d}\sigma + \int_{\Sigma_{t}} \left|\mathring{\mathrm{h}}\right|^{2} + \mathrm{R} + \frac{|\nabla^{\top} \mathrm{H}|^{2}}{\mathrm{H}^{2}} \, \mathrm{d}\sigma \right) \geq 0. \end{split}$$

 $\mathfrak{m}_H(\Sigma_t)$ is monotone <u>non-decreasing</u>.

Assume that the flow is given by the level set of a function $u: M \setminus \Omega \to \mathbb{R}$ such that $\{x: u(x) = t\} = F(t,\cdot)$. Then, for $p \in N$ we have

$$1 = \frac{\partial}{\partial t}t = \frac{\partial}{\partial t}u(F(t,p)) = \left\langle \nabla u, \partial_t F(t,p) \right\rangle = \left\langle \nabla u, \frac{\nu(F(t,p))}{\mathrm{H}(F(t,p))} \right\rangle = \frac{1}{\mathrm{H}(F(t,p))} \left\langle \nabla u, \frac{\nabla u}{|\nabla u|} \right\rangle,$$

hence,

$$H = |\nabla u|$$
.

Keeping in mind that $H = \operatorname{div}\left(\frac{\nabla u}{|\nabla u|}\right) = \Delta_1 u$, we derive the level set formulation of IMCF.

Level set formulation

Given $\Omega \,{\subseteq}\, M$ closed and bounded, we define w_1 as the solution to

$$egin{aligned} \Delta_1 w_1 &= |
abla w_1| && ext{on } M \diagdown \Omega, \ w_1 &= 0 && ext{on } \partial \Omega, \ w_1 & o +\infty && ext{as } \operatorname{d}(x,\Omega) \, o \, +\infty \end{aligned}$$

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Given $\Omega \subseteq M$ closed and bounded, we define u_2 as the solution to

$$\begin{cases} \Delta u_2 \,=\, 0 & \qquad \text{on } M {\smallsetminus} \Omega, \\ u_2 \,=\, 1 & \qquad \text{on } \partial \Omega, \\ u_2 \,\rightarrow\, 0 & \qquad \text{as } \mathrm{d}(x,\Omega) \,{\to}\, +\infty, \end{cases}$$

If a solution exists and Ω is a regular domain, then $u_2 \in C^\infty(M \setminus \mathring{\Omega})$ is unique and it is smooth till the boundary of Ω . By Sard theorem, the set of critical values of u_2 has zero Lebesgue measure. Notice that $w_2 = -\log(u_2)$ solves

$$\begin{split} \Delta w_2 &= \left| \nabla w_2 \right|^2 & \text{on } M \diagdown \Omega, \\ w_2 &= 0 & \text{on } \partial \Omega, \\ w_2 &\to +\infty & \text{as } \mathrm{d}(x,\Omega) \to +\infty \end{split}$$

Suitable variants of the Dirichlet Energy

$$\mathcal{G}_2 = \int_{\Sigma} |\nabla u|^2 \, \mathrm{d}\sigma$$

are monotone on the level set of harmonic functions.

Non-exhaustive list of applications:

- Uniqueness of smooth tangent cones at infinity in Ricci-flat manifolds [Colding 12 · Acta Math.],[Colding, Minicozzi 14 · Invent. Math.]
- Riemannian positive mass theorem [Agostiniani, Mazzieri, Oronzio '24 · CMP], [Cederbaum, León Quirós, Meco '25]
- Willmore inequality [Agostiniani, Fogagnolo, Mazzieri '20 · Invent. Math.] [Cederbaum, Miehe '25 · Trans. Am. Math. Soc.]
- Characterization of Ricci pinched 3-manifolds [Benatti, Mantegazza, Oronzio, P , '25 · J. Geom.
 An.]

2 Formal relations

 $\begin{cases} \partial_t F(p,t) = \frac{\nu(F(t,p))}{\mathrm{H}(F(t,p))} \\ F(0,p) = F_0(p) \end{cases}$

IMCF

Harmonic functions

$$\begin{cases} \Delta u_2 \,=\, 0 & \qquad \text{on } M \diagdown \Omega, \\ u_2 \,=\, 1 & \qquad \text{on } \partial \Omega, \\ u_2 \,\rightarrow\, 0 & \qquad \text{as } \mathrm{d}(x,\Omega) \,{\to}\, +\infty, \end{cases}$$

Reformulation

Level set formulation

$$\left\{egin{array}{ll} \Delta_1 w_1 &= |
abla w_1| & ext{on } M {\smallsetminus} \Omega \ & w_1 &= 0 & ext{on } \partial \Omega, \end{array}
ight.$$

on
$$\partial\Omega,$$

Change of variable

$$\begin{cases} \Delta_1 w_1 = |\nabla w_1| & \text{on } M \smallsetminus \Omega, \\ w_1 = 0 & \text{on } \partial \Omega, \\ w_1 \to +\infty & \text{as } \operatorname{d}(x,\Omega) \to +\infty, \end{cases} \qquad \begin{cases} \Delta w_2 = |\nabla w_2|^2 & \text{on } M \smallsetminus \Omega, \\ w_2 = 0 & \text{on } \partial \Omega, \\ w_2 \to +\infty & \text{as } \operatorname{d}(x,\Omega) \to +\infty, \end{cases}$$

Monotone quantities

Willmore energy
$$\mathcal{F}_1(t) = \int_{\partial \Omega_t} \!\!\!\! H^2 \, \mathrm{d}\sigma$$

Hawking mass $\mathfrak{m}_H(\Sigma) \coloneqq \sqrt{\frac{|\Sigma|}{16\pi}} \left(1 - \frac{1}{16\pi} \int_{\Sigma} \!\!\!\! H^2 \, \mathrm{d}\sigma \right)$

Dirichlet energy $\mathcal{G}_2(t) = \int_{\Sigma_1} \left| \nabla w_2 \right|^2 \mathrm{d}\sigma$

Definition - p-IMCF

Let $p \in [1,2]$. Given $\Omega \subseteq M$ closed and bounded, we define w_n as the solution to

$$\begin{cases} \Delta_p w_p &= |\nabla w_p|^p & \quad \text{on } M \diagdown \Omega, \\ w_p &= 0 & \quad \text{on } \partial \Omega, \\ w_p \to +\infty & \quad \text{as } \mathrm{d}(x,\Omega) \to +\infty, \end{cases}$$

(p-IMCF)

where $\Delta_p f = \operatorname{div}(|\nabla f|^{p-2} \nabla f)$.

We denote $\Omega_t^{(p)} = \{ w_p \leq t \}.$

- \circ p=2: linear potential theory, $u_2=\exp(-w_2)$ is harmonic.
- $\circ p \in (1,2)$: nonlinear potential theory (NPT), $u_p = \exp(-w_p/(p-1))$ is p-harmonic.
- p=1: weak inverse mean curvature flow [Huisken, Ilmanen '01 · JDG].

Proposition - [Moser '07 · JEMS], [Kotschwar, Ni '09 · Ann. Sci. Éc. Norm. Supér.]

If w_p exists for every $p \in [1,2]$, then $w_p \to w_1$ locally uniformly as $p \to 1^+$.

Aim

We want to approximate monotonic quantities along the IMCF with monotonic quantities in NPT.

Step 1: Approximation of the Willmore energy

Construction of \mathcal{F}_p proxy for $\mathcal{F}_1 = \int_{\Sigma_*} \mathbf{H}^2 d\sigma$.

Formally,

- \circ \mathcal{F}_p should coincide with \mathcal{F}_1 in the limit $p \to 1^+$.
- $\circ \ \mathcal{F}_p'$ should coincide with \mathcal{F}_1' in the limit $p \to 1^+$.
- $\circ \ \ \mathcal{F}_p \text{ should be monotone along } \Delta_p w_p = \left| \nabla w_p \right|^p.$

Consider as model case $M=\mathbb{R}^3$ and Ω with radial symmetry. Call w_p and w_1 the solutions to $\Delta_p w_p = |\nabla w_p|^p$ and $\Delta_1 w_1 = |\nabla w_1|$, respectively.

Then, $w_p=(3-p)\log|x|,\,w_1=2\log|x|$ and the relation $\frac{2}{3-p}w_p=w_1$ is satisfied. Hence

$$\frac{2}{3-p} \big| \nabla w_p \big| = |\nabla w_1| = \mathrm{H}.$$

The term $\left(\mathrm{H}-\frac{2}{3-p}|\nabla w_p|\right)^2$ is a sort of <u>deficit</u> from the model case of IMCF in \mathbb{R}^3 .

Possible candidates are

$$\mathcal{F}_p(t) = \int_{\partial \Omega_t} \mathbf{H}^2 - \left(\mathbf{H} - \frac{2}{3-p} \big| \nabla w_p \big| \right)^2 \mathrm{d}\sigma \quad \text{and} \quad \mathcal{G}_p(t) = \int_{\partial \Omega_t} \frac{4 \big| \nabla w_p \big|^2}{(3-p)^2}.$$

From direct computation we have

$$\begin{split} \mathcal{F}_{1}'(t) &= -2\int_{\partial\Omega_{t}}\frac{\left|\nabla^{\top}\mathbf{H}\right|^{2}}{\mathbf{H}^{2}} + \left|\mathring{\mathbf{h}}\right|^{2} + \mathrm{Ric}(\nu,\nu)\mathrm{d}\sigma, \\ \mathcal{F}_{p}'(t) &= -\frac{4}{3-p}\int_{\partial\Omega_{t}}\frac{\left|\nabla^{\top}|\nabla w_{p}|\right|^{2}}{\left|\nabla w_{p}\right|^{2}} + \left|\mathring{\mathbf{h}}\right|^{2} + \mathrm{Ric}(\nu,\nu) + \frac{3-p}{2(p-1)}\left(\mathbf{H} - \frac{2}{3-p}\left|\nabla w_{p}\right|\right)^{2}\mathrm{d}\sigma \\ \mathcal{G}_{p}'(t) &= \frac{4}{p-1}\int_{\partial\Omega_{t}}2\frac{\left|\nabla w_{p}\right|^{2}}{(3-p)^{2}} - \frac{\left|\nabla w_{p}\right|\mathbf{H}}{3-p}\,\mathrm{d}\sigma = \frac{1}{p-1}\left(\mathcal{G}_{p}(t) - \mathcal{F}_{p}(t)\right) \end{split}$$

If $\operatorname{Ric} \geq 0$, then $\mathcal{F}'_{p}(t) \leq 0$. Moreover,

$$\left(\mathrm{e}^{-\frac{1}{p-1}t}\mathcal{G}_p'(t)\right)' = \mathrm{e}^{-\frac{1}{p-1}t}\left(\mathcal{G}_p''(t) - \frac{1}{p-1}\mathcal{G}_p'(t)\right) = -\frac{\mathrm{e}^{-\frac{1}{p-1}t}}{p-1}\mathcal{F}_p'(t) \geq 0.$$

If, in addiction, AVR(g) > 0, this implies $\mathcal{G}_p'(t) \leq 0$, thus

$$0\!\leq\!\mathcal{G}_p(t)\!\leq\!\mathcal{F}_p(t)\!\leq\!\mathcal{F}_1(t)$$

 \mathcal{F}_p and \mathcal{G}_p

 \mathcal{F}_p is <u>monotone</u> non-increasing

- **Step 2:** "p-version" of the Hawking mass $\mathfrak{m}_H(\Sigma) \coloneqq \sqrt{\frac{|\Sigma|}{16\pi}} \left(1 \frac{1}{16\pi} \int_{\Sigma} H^2 d\sigma\right)$. We need to
- o approximate the area term, maintaining its scaling;
- o maintain the exponetial growth along the flow.

Definition - Normalized p-capacity

Let Ω be a closed, bounded set in M. The normalized p-capacity of Ω is defined as

$$\mathfrak{c}_p(\partial\Omega)\!\coloneqq\!\inf\!\left\{\frac{1}{4\pi}\!\left(\frac{p\!-\!1}{3\!-\!p}\right)^{p\!-\!1}\!\int_{M\backslash\Omega}\!|\nabla\phi|^p\!\,\mathrm{d}\mu:\phi\!\in\!\mathscr{C}_c^\infty(M),\phi\!\ge\!1\,\mathrm{on}\,\Omega\right\}$$

It holds $\mathfrak{c}_p(\partial\Omega_t) = \mathrm{e}^t\,\mathfrak{c}_p(\partial\Omega)$ and $\mathfrak{c}_p(\partial\Omega) = \int_{\partial\Omega} \left(\frac{|\nabla w_p|}{3-p}\right)^{p-1}\mathrm{d}\sigma.$

As $p \to 1^+$ we have that $\mathfrak{c}_p(\Sigma)$ tends to $|\Sigma^*|$, where Σ^* is the outward minimizing hull of Σ [Fogagnolo, Mazzieri '22 · JFA].

Natural choice:

$$\mathfrak{m}_H^{(p)}(\Sigma) \coloneqq \frac{\mathfrak{c}_p(\Sigma)^{\frac{1}{3-p}}}{2} \left(1 - \frac{1}{16\pi} \int_{\Sigma} \mathbf{H}^2 - \left(\mathbf{H} - \frac{2}{3-p} \big| \nabla w_p \big| \right)^2 \mathrm{d}\sigma \right)$$

Finally, we would like to "summarize" **setting 1** and **setting 2** in the same functional, that then can be "specialized" both to Willmore-type energies and Hawking-type masses.

$$\begin{split} \mathcal{F}_p(\partial\Omega_t) &= \mathfrak{c}_p(\partial\Omega_t)^{\frac{\alpha}{n-p}-1} \int_{\partial\Omega_t} \left|\nabla w_p\right|^{\alpha+p-2} \Big[\mathbf{H} - \Big(\frac{n-1}{n-p} - \frac{1}{\alpha}\Big) \big|\nabla w_p\big| \Big] \, \mathrm{d}\sigma \\ &+ \int_0^t \mathfrak{c}_p(\partial\Omega_s)^{\frac{\alpha}{n-p}-1} \int_{\partial\Omega_s} \left|\nabla w_p\right|^{\alpha+p-3} \mathrm{Ric}(\nu,\nu) \, \mathrm{d}\sigma \, \mathrm{d}s. \end{split}$$

With this correction term, we have a monotone functional in (M,g) noncompact, complete, Riemannian manifold of dimension $n \ge 3$, without restrictions nor on Ric neither on R of M.

IMCF

$$\begin{cases} \Delta_p w_p \,=\, |\nabla w_p|^p & \quad \text{on } M \diagdown \Omega, \\ w_p \,=\, 0 & \quad \text{on } \partial \Omega, \\ w_p \,\to\, +\infty & \quad \text{as } \mathrm{d}(x,\Omega) \,\to\, +\infty, \end{cases}$$

$$\begin{cases} \Delta_1 w_1 \,=\, |\nabla w_1| & \text{on } M \diagdown \Omega, \\ w_1 \,=\, 0 & \text{on } \partial \Omega, \\ w_1 \,\to\, +\infty & \text{as } \mathrm{d}(x,\Omega) \,\to\, +\infty, \end{cases}$$

As $p \rightarrow 1^+$, solutions converge to solutions and level sets to level sets.

$$\begin{split} \mathcal{F}_p(t) &= \int_{\partial \Omega_t} \mathbf{H}^2 - \left(\mathbf{H} - \frac{2}{3-p} |\nabla w_p|\right)^2 \mathrm{d}\sigma \\ \mathcal{F}_p'(t) &= -\frac{4}{3-p} \int_{\partial \Omega_t} \frac{|\nabla^\top |\nabla w_p||^2}{|\nabla w_p|^2} + \left|\mathring{\mathbf{h}}\right|^2 + \mathrm{Ric}(\nu, \nu) \\ &+ \frac{3-p}{2(p-1)} \left(\mathbf{H} - \frac{2}{3-p} |\nabla w_p|\right)^2 \mathrm{d}\sigma \end{split} \qquad \mathcal{F}_1(t) = -2 \int_{\partial \Omega_t} \frac{|\nabla^\top \mathbf{H}|^2}{\mathbf{H}^2} + \left|\mathring{\mathbf{h}}\right|^2 + \mathrm{Ric}(\nu, \nu) \mathrm{d}\sigma \end{split}$$

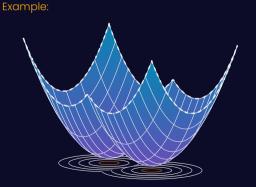
 $\mathcal{F}_p(t)$ approximate $\mathcal{F}_1(t)$, also at the level of their derivatives.

3 Limitations of the smooth setting

Theorem - [Gerhardt '90 · JDG], [Urbas '90 · Math.Z.]

Any compact, star-shaped initial surface remains star-shaped and smooth along the flow and becomes an expanding round sphere as $t \to +\infty$.

The classical IMCF could develop singularities.



The level set formulation $\Delta_1 w_1 = |\nabla w_1|$

- is a degenerate elliptic PDE,
- is not the Euler-Lagrange equation of a functional.

Existence and regularity theory for weak IMCF is delicate [Huisken, Ilmanen '01 · JDG].

p-IMCF

-IMCF

Regularity

$$w_p \in \mathcal{C}^{1,\beta} \text{ and smooth on } \left\{ \nabla w_p \neq 0 \right\}$$

$$|\nabla w_p| \in W^{1,2}$$

 \boldsymbol{w}_1 is Lipschitz

Regularity of level sets

 $|\nabla w_p|\neq 0$ almost everywhere on $\partial\Omega_t$ for a.e. t $\partial\Omega_t$ is smooth away from the critical set.

 Ω_t is strictly outward minimizing $\partial\Omega_t$ admits a weak mean curvature ${
m H}\!>\!0$ $\partial\Omega_t$ is $\mathscr{C}^{1,\beta}$ out of a set of dimension $n\!-\!8$.

Goal: Interpret the level sets as (weak) surfaces, regular enough to get a notion of mean curvature, second fundamental form, and Gauss curvature.

4 GMT meets PDE

$$\begin{split} \boldsymbol{\mathcal{F}}_{p}^{\prime}(t) = & -\frac{4}{3-p} \int_{\partial \Omega_{\ell}^{(p)}} \frac{\left| \boldsymbol{\nabla}^{\intercal} |\boldsymbol{\nabla} \boldsymbol{w}_{p}| \right|^{2}}{\left| \boldsymbol{\nabla} \boldsymbol{w}_{p} \right|^{2}} + \left| \boldsymbol{\mathring{\mathbf{h}}} \right|^{2} + \mathrm{Ric}(\boldsymbol{\nu}, \boldsymbol{\nu}) \\ & + \frac{3-p}{2(p-1)} \bigg(\boldsymbol{\mathbf{H}} - \frac{2}{3-p} |\boldsymbol{\nabla} \boldsymbol{w}_{p}| \bigg)^{2} \mathrm{d}\boldsymbol{\sigma} \end{split}$$

$$\mathcal{F}_{1}^{'}(t) = -2\int_{\partial\Omega^{(1)}} \frac{\left|\nabla^{\top} |\nabla w_{p}|\right|^{2}}{\left|\nabla w_{p}\right|^{2}} + \left|\mathring{\mathbf{h}}\right|^{2} + \mathrm{Ric}(\nu, \nu) \mathrm{d}\sigma$$

- **Goal 1:** give meaning to all the terms in the formal derivative.
- Goal 2: (strong enough) convergence result.

Theorem - [Benatti, P —, Pozzetta '24]

For every $p \in [1,2]$, almost every $\partial \Omega_t^{(p)}$ is a curvature varifold.

Moreover,

- $\circ \partial \Omega_t^{(p)}$ converges (up to subsequence) to $\partial \Omega_t^{(1)}$ in the sense of curvature varifolds for a.e. t > 0.
- $\circ w_p \to w_1 \text{ in } W_{\text{loc}}^{1,q} \text{ for } q < +\infty.$

Step 1. almost every $\partial \Omega_t^{(p)}$ is a curvature varifold.

 ε -regularization and **a** priori bounds. This gives meaning to all terms in

$$\mathcal{F}_p^{'}(t) = -\frac{4}{(3-p)} \int_{\partial\Omega^{[p)}} \frac{\left|\nabla^\top |\nabla w_p|\right|^2}{\left|\nabla w_p\right|^2} + \left|\mathring{\mathbf{h}}\right|^2 + \mathrm{Ric}(\nu,\nu) + \frac{3-p}{2(p-1)} \left(\mathbf{H} - \frac{2}{3-p} \left|\nabla w_p\right|\right)^2 \mathrm{d}\sigma.$$

Step 2. $\partial \Omega_t^{(p)}$ converges to $\partial \Omega_t^{(1)}$ in the sense of curvature varifolds for a.e. t > 0.

- $1 |\partial \Omega^{(p)}| \xrightarrow{p \to 1^+} |\partial \Omega^{(1)}|$ for almost every t,
- $2. \int_{\partial\Omega^{(p)}} (|\nabla w_p| \mathbf{H})^2 d\sigma \xrightarrow{p \to 1^+} 0.$
- 3. $\int_{\partial\Omega^{(p)}} |\mathbf{h}|^2 d\sigma$ is bounded.

Step 3. Gradient convergence.

Combine: $\circ \nabla w_n \xrightarrow{p \to 1^+} \nabla w_1$ weakly in L^q

$$ullet \left\|w_p
ight\|_{L^1_{ ext{loc}}} \stackrel{p o 1^+}{\longrightarrow} \left\|w_1
ight\|_{L^1_{ ext{loc}}}$$

$$\circ \lim_{p \to 1^+} \int_0^T \int_{\partial \Omega^{(p)}} |\nu^{(p)} - \nu^{(1)}|^2 d\sigma dt = 0$$

1. + 2. \Rightarrow convergence in the sense of varifolds. 2. + bound on $\int |\mathring{\mathbf{h}}|^2 \Rightarrow 3$.

1. $+ 2. + 3. \Rightarrow$ convergence as curvature varifolds.

For $p \in (1,2)$, we do not have enough regularity to define an Euler characteristic that encompass the topological properties of the level sets.

A weak induced scalar curvature can be defined through an integral version of the Gauss equation:

$$\int_{\Sigma} \mathbf{R}^{\top} \mathbf{d}\sigma = \int_{\Sigma} \mathbf{R} - 2 \mathrm{Ric}(\nu, \nu) + \mathbf{H}^2 - |\mathbf{h}|^2 \, \mathbf{d}\sigma$$

Theorem - [Benatti, P —, Pozzetta '24]

Let (M,g) be a complete, noncompact 3-dimensional Riemannian manifold. Let $p\in (1,2)$ and w_p be the solution to $\Delta_p w_p = |\nabla w_p|^p$. Then, for almost every $t\in [0,+\infty)$ it holds

$$\int_{\partial\Omega^p} \mathbf{R}^\top \mathbf{d}\sigma \in 8\pi \mathbb{Z}.$$